

Top++: a program for the calculation of the top-pair cross-section at hadron colliders

Michał Czakon^a, Alexander Mitov^b

^a*Institut für Theoretische Teilchenphysik und Kosmologie, RWTH Aachen University, D-52056 Aachen, Germany*

^b*Theory Division, CERN, CH-1211 Geneva 23, Switzerland*

Abstract

We present the program **Top++** for the numerical evaluation of the total inclusive cross-section for producing top quark pairs at hadron colliders. The program calculates the cross-section in a) fixed order approach with exact NNLO for $q\bar{q} \rightarrow t\bar{t} + X$, $q q \rightarrow t\bar{t} + X$, $q\bar{q}' \rightarrow t\bar{t} + X$, $q q' \rightarrow t\bar{t} + X$ and through approximate NNLO for $g g \rightarrow t\bar{t} + X$ and b) by including soft-gluon resummation for the hadronic cross-section in Mellin space with full next-to-next-to-leading logarithmic accuracy. The program offers the user significant flexibility through the large number (31) of available options. **Top++** is written in C++. It has a very simple to use interface that is intuitive and directly reflects the physics. The running of the program requires no programming experience from the user.

Program summary

<i>Name of the program:</i>	Top++ (ver. 1.3.1).
<i>Program's homepage:</i>	http://www.alexandermitov.com/software
<i>License, Warranty:</i>	GNU Public License. No warranty given or implied.
<i>Compiler:</i>	Developed and tested with GNU Compiler Collection's C++ compiler.
<i>Operating system:</i>	Linux; Mac OS X; can be adapted for other unix systems.
<i>Program language:</i>	C++.
<i>Memory required to execute:</i>	Typically less than 200 MB.
<i>External libraries:</i>	GNU Scientific Library (GSL); the Les Houches Accord PDF Interface (LHAPDF).
<i>Keywords:</i>	Top-quark, Resummation, QCD, Precision Physics, Hadron Colliders.
<i>Typical running time:</i>	Depending on the options. The program is optimized for speed.
<i>Accuracy:</i>	Sub per-mill accuracy achievable in realistic time (program does not employ Monte Carlo methods).

*Preprint numbers: CERN-PH-TH/2011-303, TTK-11-58

What is new in ver.1.3.1: changes from ver.1.3.

- Bug fix (in file `top++.cpp`): the program could call the wrong prescription for the calculation of pdf uncertainties with the NNPDF family of sets. The bug does not affect earlier versions of `Top++` or any other functionality. In particular, all numbers in this manual are unaffected. We thank Sebastian Naumann-Emme for bringing it up to our attention.

What is new in ver.1.3: changes from ver.1.2.

- Added the complete NNLO results [1] for the reactions:

$$q\bar{q} \rightarrow t\bar{t} + q\bar{q}|_{\text{NS}}, \quad (1)$$

$$q\bar{q}' \rightarrow t\bar{t} + q\bar{q}', \quad (2)$$

$$qq' \rightarrow t\bar{t} + qq', \quad (3)$$

$$qq \rightarrow t\bar{t} + qq. \quad (4)$$

Adding the reaction $q\bar{q} \rightarrow t\bar{t} + q\bar{q}|_{\text{NS}}$ (as defined in Ref. [1]) completes the calculation of the reaction $q\bar{q} \rightarrow t\bar{t} + X$ at NNLO. Please note that, compared to version 1.2, version 1.3 adds few seconds to each evaluation due to the initialization of the newly added partonic fluxes $qq, qq', q\bar{q}'$.

- Added an option that allows switching between “LHgrid” and “LHpdf” files in the LHAPDF interface. The option is called `PdfFileType` with default value `LHgrid`. The alternative value is `LHpdf`.
- Added the symmetric prescription for computing PDF uncertainties (defined as in Ref. [2]). It is automatically invoked for the ABM11 [2], ABKM09 [3] and earlier Alekhin [4] family of PDF sets.

What is new in ver.1.2: changes from ver.1.1.

With respect to version 1.1 (minor bug fixes with respect to version 1.0), `Top++1.2` includes the exact NNLO result for $q\bar{q} \rightarrow t\bar{t} + X$. This is a major upgrade to the program’s functionality that necessitates the following modifications:

- Remove the option `Cbarqq` which is now obsolete due to the inclusion of the exact NNLO $q\bar{q} \rightarrow t\bar{t} + X$ result.
- The functionalities of the following options are modified with respect to ver.1.1 and lower: `OrderF0`, `OrderRES`, `RESON0Roff`, `H2qq`, `H2gg1`, `H2gg8`, `NNLO`, `NNLOon0Roff`, `Cbargg`, `ECMLHC`.
- The description of the option `MtopLimit` is improved.
- Included is a new set of examples, that reflects the new functionality.
- Slight modifications to the Manual that reflect these changes.
- Added functionality to the file `top++.cpp` that restricts the option `MtopStep` to be positive (see its description below), and prevents the program from possibly entering into an infinite loop.

The default settings of the program reproduce the Tevatron NNLO numbers of Ref. [5]. The best recommended setting for the LHC can be found in the file `top++best-LHC8.cfg`. They are a combination of the findings of Ref. [6] with the exact NNLO result for the $q\bar{q} \rightarrow t\bar{t} + X$ reaction from Ref. [5]. For the sake

of documenting the results,¹ in the following we explicitly present our best predictions for the Tevatron, LHC @ 7 TeV and LHC @ 8 TeV:

$$\sigma_{\text{tot}}^{\text{NNLO+NNLL}}(\text{Tevatron}; m_t = 173.3 \text{ GeV}) = 7.067^{+0.143(2.0\%)}_{-0.232(3.3\%)} [\text{scales}]^{+0.186(2.6\%)}_{-0.122(1.7\%)} [\text{pdf}], \quad (5)$$

$$\sigma_{\text{tot}}^{\text{NNLO+NNLL}}(\text{LHC}_{7\text{TeV}}; m_t = 173.3 \text{ GeV}) = 156.2^{+14.6(9.3\%)}_{-7.9(5.0\%)} [\text{scales}]^{+4.1(2.6\%)}_{-4.2(2.7\%)} [\text{pdf}], \quad (6)$$

$$\sigma_{\text{tot}}^{\text{NNLO+NNLL}}(\text{LHC}_{8\text{TeV}}; m_t = 173.3 \text{ GeV}) = 222.5^{+21.0(9.4\%)}_{-11.4(5.1\%)} [\text{scales}]^{+5.4(2.4\%)}_{-5.6(2.5\%)} [\text{pdf}]. \quad (7)$$

The numbers for the Tevatron Eq. (5) reproduce the best prediction of Ref. [5], while the ones for the LHC Eqs. (6,7) are new, and represent the current state-of-the-art.

The dramatic improvement from the inclusion of the NNLO corrections to $q\bar{q} \rightarrow t\bar{t} + X$ at the Tevatron has already been discussed in Ref. [5]. The inclusion of the NNLO correction in $q\bar{q} \rightarrow t\bar{t} + X$ improves the theoretical prediction at the LHC with respect to the NLO+NNLL one [6]: for both LHC @ 7 TeV and LHC @ 8 TeV the scale uncertainty bands perfectly overlap, and the total scale uncertainty in Eqs. (6,7) is reduced, in absolute size, by about 2%.

What is new in ver.1.1: changes from ver.1.0.

With respect to version 1.0, **Top++1.1** fixes two bugs in the file `Resummation.cpp`.

- Corrects the normalization of the hard matching constants `H2qq`, `H2gg1`, `H2gg8` (the factor of $1/\pi^2$ was missing in ver.1.0).
- Fixes a bug in the calculation of the matching of the all-order soft exponent. The bug was affecting the results in the unusual cases when lower logarithmic accuracy was requested than the fixed order one, i.e. the following three cases: NLO+LL, NNLO+LL and NNLO+NLL.

These bug fixes do not affect any of the numbers published in this manual or in Ref. [6].

1. Introduction: what is this program for?

This program calculates the top-pair total inclusive cross-section in hadron collisions. It has built-in the most advanced theoretical results available in the literature. The program is able to calculate the total top-pair cross-section in both pure fixed order perturbation theory and by including soft-gluon resummation through next-to-next-to-leading logarithmic order (NNLL). **Top++** is the first publicly available program that can perform soft gluon resummation in top-pair production.

The program is written in C++ in a modern, modular and object-oriented way. It should be very easy to install on most Linux systems; please consult Appendix A for details. Moreover, once installed, the program is trivial to run. The program has been written with a user in mind that has no programming experience whatsoever. For that reason, the program has a very simple user interface that is the only point of contact between the program and the user. The user interface is described in Section 2. Users with average programming experience will find the program very easy to customize to their own needs. That may not be needed, however, since all options that have arisen in our own work on the subject have already been pre-programmed and are easy to access directly through the user interface.

In this manual we do not describe the physics in detail. A short description, needed to make the reading of the program's options self-consistent, can be found in Section 4. The relevant description of the physics

¹We take this opportunity to update the legacy NLO+NLL numbers given in Tables 2,3 of Ref. [6]. The fully updated results, which can be obtained with any version of **Top++**, read: $172.0^{+15.4(8.9\%)}_{-13.7(7.9\%)}$ (Table 2, line 2); $165.8^{+13.7(8.3\%)}_{-14.1(8.5\%)}$ (Table 2, line 3); $7.039^{+0.197(2.8\%)}_{-0.422(6.0\%)}$ (Table 1, line 2); $6.911^{+0.269(3.9\%)}_{-0.497(7.2\%)}$ (Table 1, line 3). We thank Christian Schwinn for bringing the need for this update to our attention.

and the options implemented in this program can be found in Refs. [5, 6]. The results in these papers represent the most advanced results in top physics to date. Given the fast pace of developments in the subject, however, we expect that the complete set of NNLO results will be available soon.

And an important word of caution regarding the interpretation of numerical results derived with the help of the program. Clearly a program with a number of options (like ours) is capable of correctly producing almost any number. To be able to interpret such a number, however, one needs to know exactly how it was derived. This can become a problem with the state of top physics where the most advanced results used to constitute a series of approximations (see Ref. [6]) for both Tevatron and LHC. To reduce the chance of mis-communicating results derived with this program, we have implemented default settings that correspond to the best prediction of Ref. [5] for the Tevatron and a combination of Refs. [5, 6] for the LHC, and we urge the users to use them for publications. Parameters that are not considered part of these settings are the parton distribution functions (PDF) and the value of the top quark mass. These should always be quoted to ensure complete documentation of the results. Other settings, mainly related to the workings of the program, can affect the final numerical value (through numerical precision) but that should be at a precision level below what is physically relevant. As a rule of thumb we urge the user to produce numbers that are correct at the level of one per-mill (i.e. $\mathcal{O}(10^{-3})$). Precision better than that is a pure matter of taste on the side of the user (and, at a certain level, of the ability of the otherwise very capable integration routines).

2. How to use the program

In the following we assume that the program has been installed and is running correctly; see Appendix A for more details. Next, we assume that the program is installed in a directory called `top++` and we are already there, i.e. the command line reads: ²

```
~ top++$
```

Next one needs to open the file `top++.cfg` in a text editor. Using, for example, *pico* one types:

```
~ top++$ pico top++.cfg
```

Once the user has set all options at their desired values, the file needs to be saved. A number of examples are supplied with the program; see Appendix B. Then the user needs to execute the program by typing:

```
~ top++$ ./top++
```

The program starts running, displays its step-by-step progress, timing pre-defined milestones. Once the run is completed, the final result is displayed and the program exits. The output can be found on the screen and in the file `top++.res` located in the program's directory. If a new run is desired, one has to simply repeat all the steps described above.

In the following we describe all 30 options that are available to the user through the file `top++.cfg`. The options are grouped into five subgroups. Please note that since all options have predefined default values, the user needs to only specify the values for options that differ from their default values. In particular, if all options are at their default values the file `top++.cfg` can be empty. It might be convenient to keep certain options typed in the file, but when not in use the user can comment them out by putting the symbol `/` at the beginning of a line.

1. **General Setup** (type of collider, PDF set, pure fixed order calculation versus one with resummation).

- (a) **Collider**: Takes two values: `TEV` (default) or `LHC` as labels for $\bar{p}p$ or pp colliders.

²Please note that the precise text on the left of the cursor depends on the particular terminal.

- (b) **WithResummation**: If set to **YES** (default), the program will compute the observable at $N^n LO + N^k LL$ depending on the values for n and k specified in the subgroup **Resummation** below. In this case all settings in the subgroup **Fixed Order** become irrelevant.
- (c) **PDFuncertainty**: If set to **YES** the program will compute and display the PDF uncertainty. Three methods for computing PDF uncertainties are available: First, the symmetric prescription (as defined in [2]) is automatically invoked for the ABM11 [2], ABKM09 [3] and earlier Alekhin [4] family of pdf sets. Second, the prescription for the NNPDF family of sets [7] is selected automatically for any NNPDF set and, third, for all other PDF sets the program uses by default the asymmetric prescription of Ref. [8] (see also Ref. [9]). In particular, the latter prescription is the most appropriate one for the MSTW2008 family of sets [10]. When the PDF uncertainty is being calculated, central scales choice $\mu_F = \mu_R = m_t$ is set automatically, independently of the values for these scales set by the user (as described below). If **PDFuncertainty** is set to **NO** (default) then the program computes and displays the scale variation for the choice of μ_F, μ_R specified by the user (explained below) for a single pdf member (specified by the user; see option 1f below). Please note that the definition of a “central” set depends on the PDF set.
- (d) **RestrictedScaleVariation**: Takes either **NO** or a number greater than or equal to 1.0. This option allows the user to set any restriction (or no restriction at all if **RestrictedScaleVariation** **NO**) on the allowed ratio for the renormalization and factorization scales. For more information see the description of the group of options **Top quark mass and renormalization/factorization scales**. The default value is 2.0 which corresponds to the restricted scale variation of Refs. [5, 6].
- (e) **PDFset**: The PDF set. The program uses the LHAPDF library and follows its nomenclature. The default set is `MSTW2008nnlo68c1`. Both “.LHgrid” and “.LHpdf” files can be used, see the option **PdfFileType** below.
- (f) **PDFmember**: The specific member of the PDF set **PDFset** the user wishes to use for the calculation of the cross-section. Please note that the definition of a “central” set depends on the PDF set. Default value is 0.

2. **Top quark mass and renormalization/factorization scales** Please note that the ranges of μ_R and μ_F (defined in the following) that the user may provide need not be equal in length. The range of their ratio is unrestricted if **RestrictedScaleVariation** **NO** is chosen. When the option **RestrictedScaleVariation** takes a numerical value (with the constraint **RestrictedScaleVariation** ≥ 1.0), then the ratio of the two scales is restricted between:

$$\frac{1}{\text{RestrictedScaleVariation}} \leq \mu_F/\mu_R \leq \text{RestrictedScaleVariation}.$$

- (a) **Mtop**: The value of the on-shell top mass (in GeV). If the user would like to loop over a range of values for the top mass (see the description of the following two options) **Mtop** represents the lower end of that range. Default value is 173.3.
- (b) **MtopLimit**: The upper limit of the range of values for the top mass that the user wants to loop over. If **MtopLimit** $<$ **Mtop** then the program automatically sets **MtopLimit** = **Mtop** and issues a warning message. The default is the unattainable value **MtopLimit** = -1. The value -1 is converted internally to **MtopLimit** = **Mtop** too, but is special, because only for it the (annoying) warning message is suppressed.
- (c) **MtopStep**: A positive number (need not be an integer) that specifies the step with which the value of the top mass is incremented in a loop. Its default value is 1.

- (d) **muR**: Set of values for the renormalization scale (in units of the top mass **Mtop**). By default it takes the value 1.0. It can take an arbitrary number of values. In particular, to determine the scale variation of a cross-section, one may use the set:

muR 0.5 1.0 2.0 .

A fine scan can be achieved with the help of the following set:

muR 0.50 0.55 0.60 0.65 0.70 0.75 0.80 0.85 0.90 0.95 1.0 1.1 1.2 1.3 1.4 1.5 1.6 1.7 1.8 1.9 2.0 .

The larger set takes much longer to compute and typically returns the same result as the restricted set consisting of three elements. We have encountered, however, exceptional situations that in our experience lead to differences of up to 0.5%.

- (e) **muF**: Set of values for the factorization scale (in units of the top mass **Mtop**). It is set and used independently of the renormalization scale described above. Its default value is 1.0 and can be set analogously to **muR**.

3. **Resummation** (all options in this group are irrelevant when **WithResummation** **NO** is chosen).

- (a) **OrderFO**: Takes values **L0**, **NLO** or **NNLO** (default). It specifies the order n of the fixed order cross-section that should be added to the resummed result, i.e. $N^n LO + N^k LL$. For example **L0** implies $n = 0$, **NLO** implies $n = 1$, etc. The option **NNLO** includes the exactly known NNLO result for $q\bar{q} \rightarrow t\bar{t} + X$ and the approximately known NNLO cross-section [11] for $gg \rightarrow t\bar{t} + X$; see also the group of options **Fixed Order** below.
- (b) **OrderRES**: Takes the values **LL**, **NLL** and **NNLL** (default). It represents the logarithmic order k of the resummation, i.e. $N^n LO + N^k LL$. For example **LL** implies $k = 0$, **NLL** implies $k = 1$, etc. Full **NNLL** resummation is implemented matched to a hard function with **H2qq** = 84.81 [5], and the still unknown coefficients **H2gg1**, **H2gg8** set to zero by default (see also the description of the options below).
- (c) **A**: The value of the parameter A introduced in Ref. [12]. Typically used with the values **A 2** (default) and **A 0**.
- (d) **RESonORoff**: Takes the values **ON** or **OFF** (default); specifies how the unknown two-loop hard function coefficient **H2gg1**, **H2gg8** are treated. It does not affect **H2qq**. We have implemented this option consistently both for the fixed order (relevant if **OrderFO** is set to **NNLO**) and the resummed components. Moreover, the unknown constant C_{gg} in the fixed order NNLO approximation is automatically expressed through **H2gg1**, **H2gg8** (described in option 3f).
- (e) **TwoLoopCoulombs**: Takes the values **YES** (default) and **NO**. It includes/excludes the two loop Coulombic terms in the function $\sigma^{(Coul)}$, see Eq. (9) below.
- (f) **H2qq**, **H2gg1**, **H2gg8**: The three constants in the two-loop hard function. They are defined in a normalization α_S/π . **H2gg1**, **H2gg8** are irrelevant when **RESonORoff** is set to **OFF**. When **RESonORoff** is set to **ON**, one needs to vary them to determine the full theoretical uncertainty. The default values for **H2gg1** and **H2gg8** are 0.0 and for **H2qq** is 84.81 [5].

4. **Fixed Order** (all options in this group are irrelevant when **WithResummation** **YES** is chosen).

- (a) **L0**: takes values **YES** (default) or **NO**.
- (b) **NLO**: takes values **YES** (default) or **NO**. For example, to perform a calculation in NLO QCD one needs to set **L0 YES** and **NLO YES**. This way the user has direct access to the individual orders in the perturbative expansion. Same applies to the option **NNLO** below.

- (c) **NNLO**: takes values **YES** (default) or **NO**. This option implements the exact NNLO result for $q\bar{q} \rightarrow t\bar{t} + X$ and the threshold approximation to the NNLO cross-section for $gg \rightarrow t\bar{t} + X$ [11].
- (d) **NNLOonORoff**: Takes the values **ON** or **OFF** (default). Specifies how the $\ln(\mu)$ terms in the approximate NNLO result in $gg \rightarrow t\bar{t} + X$ are treated; see Ref. [6]. Irrelevant for $q\bar{q} \rightarrow t\bar{t} + X$.
- (e) **Cbargg**: The value for the yet-unknown two-loop constant C_{gg} in the threshold expansion of the NNLO $gg \rightarrow t\bar{t} + X$ cross-section [11] (defined in a normalization $\alpha_S/(4\pi)$). The default value for this constant, to be used at the Tevatron, is **Cbargg** = 0.0. As explained in Ref. [6] a more appropriate value for the LHC is **Cbargg** = -1136.81 (corresponding to **H2gg1** = **H2gg3** = 0).

5. **Setup parameters** (parameters related to the working of the program and other, less frequently modifiable parameters).

- (a) **ECMLHC**: The c.m. energy of the pp collider (in GeV). The default value is 8000.
- (b) **ECMTEV**: The c.m. energy of the $\bar{p}p$ collider (in GeV). The default value is 1960.
- (c) **Precision**: Defines the required relative precision of the integration routines:

$$\text{Relative precision} = 10^{-\text{Precision}}.$$

The default value **precision 2** tends to produce fast and accurate (at the per-mill level) results.

- (d) **NPdfGrid**: Defines the size of the grid on which the PDF fluxes are being discretized. The default value **NPdfGrid 100** tends to produce fast and sufficiently accurate results.
- (e) **ETA**: Parameter that controls the subtraction flux implemented in the resummed calculation. The default value **ETA 1e-5** is optimal.³ The user normally will not need to be concerned with this parameter.
- (f) **CMP**: The value where the contour for the Mellin inverse transform crosses the real line. The default value **CMP 2.7** is optimal. The user normally need not be concerned with this parameter.
- (g) **PdfFileType**: Option for switching between “LHgrid” and “LHpdf” files in the LHAPDF interface. The default value is **LHgrid**. The alternative value is **LHpdf**.

3. Once the program is running: some fine tuning

3.1. Numerical precision and speed

As every program for numerical calculations, **Top++** has its limits, too. In the following we discuss this, as well as ways to improve the accuracy and length of the runs.

There are two places where speed (and therefore accuracy) can be controlled. The first one is through the option **Precision** of the integration routines. In our own experience the value **Precision 2** is more than adequate to calculate the $t\bar{t}$ cross-section to per-mill accuracy. Increasing the value of **Precision** slows the calculation down. The program is set in such a way, that if the integration routines cannot reach the accuracy in a point, a warning message is displayed specifying the relative error returned by the integration routines. The presence of these messages, as such, is harmless. They can be ignored if the displayed relative precision is high enough. Rarely, the output might contain the symbol NaN, which stands for Not-a-Number. NaNs are returned due to invalid numerical operations such as division by zero, or operations, which produce

³We use the usual notation: $1e-p \equiv 10^{-p}$.

numbers out of the range of double precision. This is a sign of numerical instability. We have taken the pragmatic approach of not handling such exceptions in any specific way (which is very difficult in practice). In case of a NaN result, the user should rerun the calculation with increased requested accuracy.

We recommend that at least once the user does the calculation with a larger value for `Precision` and verifies that the change in the result is beyond the required accuracy.

A second (and independent) source of numerical uncertainty is the size `NPdfGrid` of the grid over which the partonic fluxes are approximated. We have implemented a second order finite difference scheme. The relative precision scales as $\sim 1/\text{NPdfGrid}^2$. In practice we have noticed that a value `NPdfGrid` 100 is more than adequate in terms of accuracy and produces very fast calculations. We recommend that at least once the user does the calculation with a larger value, say `NPdfGrid` 500, and verifies that the change in the result is beyond the required accuracy.

Finally, the user should keep in mind that the overall numerical accuracy is a combination of the settings for `Precision` and `NPdfGrid`, i.e. increasing only one of them to an extreme may not have a net positive effect on the overall uncertainty but might lead to a significant slowdown.

3.2. Outputting the results

The program outputs on the screen the information about the timing of each step and all results. The specifics depend on the requested options. At the end of the calculation, a summary of the final result is displayed, including the central value (i.e. the value corresponding to central scale choice, if requested) and the scale or PDF variation's absolute and relative values. The program also outputs the final result in a file `top++.res` which is ready for plotting with the program *Gnuplot*. This is particularly useful for the case when the user requests a loop over a range of values of the top quark mass. In such case the result for each value is conveniently recorded in the file. Please note that the file `top++.res` is overwritten after each run.

3.3. Modifying the default values of the parameters

The default values for all 30 parameters available to the user are set inside the function `main()` located in the file `top++.cpp`. All default settings can be modified by the user, although this is not recommended and should not be necessary.

3.4. Additional

The strong coupling constant is calculated at a scale μ_R through the LHAPDF interface [13]. The calculation of the cross-section is performed in a scheme with $N_F = 5$ active flavors. For consistency only PDF sets with $N_F = 5$ active flavors should be used for scales above m_{top} .

4. Contact with physics

The program `Top++` comes with settings that reproduce the “best” numbers from Ref. [5] for the Tevatron. The best numbers for the LHC will now slightly differ from the ones in Ref. [6] due to the improved precision after the inclusion of the NNLO result for $q\bar{q} \rightarrow t\bar{t} + X$. These are the best current theoretical predictions with realistic theoretical uncertainties. All physics-related inputs and approximations are defined and explained in Ref. [6], which should be consulted if more detailed questions arise.

The program calculates the $t\bar{t}$ total inclusive cross-section, schematically given as:

$$\sigma_{\text{tot}} = \sigma_{\text{F.O.}}^{(n)} + [\sigma^{\text{res}} - \sigma^{\text{res}}|_{\alpha_S^n}] . \quad (8)$$

The power n is set by the option `OrderFO` if resummation is desired (i.e. when `WithResummation YES`). The convention is that $n = 2$ corresponds to LO, $n = 3$ to NLO and $n = 4$ corresponds to NNLO. If `WithResummation NO`, i.e. one wishes to calculate in a fixed order approach, then the terms in the square bracket in Eq. (8) are absent and n is set through the options `LO`, `NLO` and `NNLO`.

In Mellin N -space, the resummed partonic cross-section reads

$$\sigma_{N,\mathbf{I}}^{\text{res part}} = \sum_{\mathbf{I}=1,8} \sigma_{N,\mathbf{I}}^{(\text{Coul})} \times \sigma_{N,\mathbf{I}}^{(\text{Hard})} \times \Delta_{N,\mathbf{I}}. \quad (9)$$

The function $\sigma^{(\text{Coul})}$ contains the Coulombic effects and has a known perturbative expansion through NNLO. The function $\sigma^{(\text{Hard})}$ is an N -independent hard function. The NNLO constant terms `H2gg1`, `H2gg8` are currently unknown. These constants, in an expansion in α_S/π , correspond to the options `H2gg1`, `H2gg8`. The option `H2qq` is available to the user, too. However, since it is set by default to the exact value derived in Ref. [5], normally it should not be modified. The $\ln(\mu)$ terms in the two-loop gg -contribution to the function $\sigma^{(\text{Hard})}$ are included if `RESonORoff ON` and excluded when `RESonORoff OFF`. The latter option gives a more conservative (i.e. realistic) determination of the theoretical uncertainties through scale variation alone.

The function $\Delta_{N,\mathbf{I}}$ contains the towers of LL, NLL and NNLL soft $\ln(N)$ logs.

By setting the option `OrderRES LL` the program included only the LL corrections from $\Delta_{N,\mathbf{I}}$, as well as the LO terms in the functions $\sigma^{(\text{Coul})}$ and $\sigma^{(\text{Hard})}$. If `OrderRES NLL` then $\Delta_{N,\mathbf{I}}$ includes LL and NLL logs, while the functions $\sigma^{(\text{Coul})}$ and $\sigma^{(\text{Hard})}$ are given by their NLO expansions. Similarly, for `OrderRES NNLL` $\Delta_{N,\mathbf{I}}$ includes LL, NLL and NNLL logs, and the functions $\sigma^{(\text{Coul})}$ and $\sigma^{(\text{Hard})}$ are given by their expansions through NNLO. The expressions for these functions can be found in Refs. [6, 5].

For pure fixed order calculations one also has a variety of options for the treatment of the currently approximately known NNLO corrections in $gg \rightarrow t\bar{t} + X$. Setting `NNLOonORoff ON` one includes the two-loop $\ln(\mu)$ terms unaccompanied by a power of β or $\ln\beta$. By setting `NNLOonORoff OFF` one excludes these two-loop logarithmic terms which corresponds to the standard setting adopted in Ref. [6]. Note that the $q\bar{q} \rightarrow t\bar{t} + X$ reaction is unaffected by this option.

Through the options `Cbargg` one has a direct access to the currently unknown two-loop constant $C_{gg}^{(2)}$. This is helpful for determining the sensitivity of the approximate NNLO cross-section to this unknown parameter.

5. Summary

We present the C++ program `Top++` for the calculation of the total top-pair cross-section at hadron colliders. This is the first publicly available program capable of performing soft-gluon resummation for this collider observable. The program uses all currently available theoretical results. The user has access to 30 options which results in a great deal of flexibility and control over the calculation on the side of the user. In this manual we have given only a very short introduction to the physics behind our program. For further details the user should consult Ref. [6].

The program is organized in a modular, object oriented way. It can seamlessly accommodate future results, notably the remaining NNLO results, by combining them with the currently available NNLL resummation and known NNLO results for $q\bar{q} \rightarrow t\bar{t} + X$ [5].

The program is optimized for speed given the user's requirements for accuracy. Our experience shows that, depending on the chosen options, the run times can vary significantly. In practice the speed can be an issue only for resummed calculations, due to the integration of rapidly oscillating functions in the complex plane. For fixed order calculations the run times are very short.

As we already emphasized `Top++` is the only publicly available program capable of producing resummed calculations. Another program for the calculation of the total cross-section in fixed order perturbation theory is available in the literature [14]. While we have not performed detailed side-by-side comparisons, our experience shows that our program can be significantly faster, especially when higher accuracy is required (since we do not use Monte Carlo integration methods). We would like to also note that a direct comparison between the results of `Top++` and `HATHOR` [14] for the approximate NNLO results may not be possible. The reason is that different implementations of this approximation are contained in both programs. In `Top++` we have implemented two versions of this approximation, depending on whether the two-loop scale dependent terms are included or excluded (see Ref. [6] for details). This is controlled through the option `NNLOonORoff OFF/ON`. An approximation analogous to our `NNLOonORoff ON` is implemented in `HATHOR` by

using the option `LOG_ONLY`. We note that while the two are compatible they are not identical, since different subleading terms $\mathcal{O}(\beta)$ are included in both implementations.

We would also like to stress (see Ref. [6]) that approximations like `NNLOonORoff ON` could significantly underestimate the theoretical uncertainty if its estimate is based on scale variation alone.

Acknowledgments

We thank Matteo Cacciari, Michelangelo Mangano and Paolo Nason for numerous cross-checks and a very fruitful collaboration [6] that led to the creation of this software. We also thank Juan Rojo and Malgorzata Worek for useful suggestions. The work of M.C. was supported by the Heisenberg and by the Gottfried Wilhelm Leibniz programmes of the Deutsche Forschungsgemeinschaft, and by the DFG Sonderforschungsbereich/Transregio 9 “Computergestützte Theoretische Teilchenphysik”.

Appendix A. Installation

The program has been written in standard C++ and has been tested to correctly compile under the GNU compiler g++ from version 4 upwards. It requires two external libraries:

- **GNU Scientific library**, which can be downloaded free of charge from <http://www.gnu.org/s/gsl/> and is used for special functions and integration
- **the Les Houches Accord PDF Interface**, which can be downloaded free of charge from <http://projects.hepforge.org/lhapdf/> and is used for the parton distribution functions

To setup the program for installation, it is necessary to set three variables in the `Makefile` contained in the `Top++` installation directory.

- `CXX` - c++ compiler. specify the full path if necessary
- `GSLDIR` - prefix directory for the gsl library, it is assumed that the library is in `GSLDIR/lib` and the include files in `GSLDIR/include`
- `LHADIR` - prefix directory for the Les Houches PDF library, it is assumed that the library is in `LHADIR/lib` and the include files are in `LHADIR/include/LHAPDF`

In the next step, it is sufficient to compile the code using

```
~ top++$ make
```

The program should compile without any error messages or warnings, and is ready to use under the name `top++`.

Alternatively to modifying the `Makefile` it is possible to compile directly with

```
~ top++$ make CXX="user CXX value" GSLDIR="user GSLDIR value" LHADIR="user LHADIR value"
```

The quotation marks above are not necessary, unless the paths contain spaces or other special values (as usual under unix).

Appendix B. Examples (updated for ver.1.3 and ver.1.3.1)

A number of examples can be found in the directory `examples`. The user can copy/paste their content into the file `top++.cfg` located in the program's main directory.

Example 1: If one executes the program with its default settings (provided in the file `top++default.cfg`) one should get the following result (this result can also be obtained with an empty file `top++.cfg` or, alternatively, one with all its options commented out):

$$\text{sigma_tot} = 7.06698 \text{ [pb]}.$$

Example 2: With the file `top++best-TEV.cfg` one can compute the scale variation of the cross-section corresponding to the best prediction of Ref. [5] for the Tevatron:

$$\text{sigma_tot} = 7.06698 + 0.143258(2.02715\%) - 0.231754(3.27939\%) \text{ [pb]}.$$

Example 3: We recommend the setting in the file `top++best-LHC8.cfg` for the computation of the cross-section for LHC @ 8 TeV (or any other LHC energy after adjusting the value of `ECMLHC`). It is based on the results of Refs. [5, 6] and reads:

$$\text{sigma_tot} = 222.664 + 21.2071(9.52428\%) - 11.3969(5.11845\%) \text{ [pb]}.$$

This result perfectly overlaps with the LHC @ 8 TeV number quoted in Eq. (26) of Ref. [6] and, as anticipated, has about 2% smaller total scale uncertainty.

Example 4: To derive the PDF variation for the best prediction of Refs. [5] at the Tevatron one can use the file `top++pdfvar-TEV.cfg`:

$$\text{sigma_tot} = 7.06698 + 0.186404(2.63768\%) - 0.122316(1.73081\%) \text{ [pb]}.$$

Example 5: To derive the PDF variation for the best prediction of Refs. [5] at the LHC @ 7 TeV one can use the file `top++pdfvar-LHC7.cfg`:

$$\text{sigma_tot} = 156.272 + 4.12863(2.64196\%) - 4.19779(2.68622\%) \text{ [pb]}.$$

Example 6: Fixed order calculations can be done by turning off the option for resummed calculation. For example, the scale variation of the NLO result at the Tevatron, computed for consistency with an NLO PDF set (MSTW in this case), is obtained from the file `top++NLO-tev.cfg`:

$$\text{sigma_tot} = 6.68123 + 0.35577(5.32492\%) - 0.752327(11.2603\%) \text{ [pb]}.$$

Example 7: The fixed order, partially NNLO result for the LHC @ 7 TeV, with excluded two-loop scale dependent terms in $gg \rightarrow t\bar{t} + X$ [6] and non-vanishing two-loop constant $C_{gg}^{(2)}$, is obtained from the file `top++Fixed-Order-LHC7.cfg`:

$$\text{sigma_tot} = 154.779 + 11.0044(7.10975\%) - 8.59923(5.55581\%) \text{ [pb]}.$$

This result overlaps very well with the corresponding number quoted in Table 2 of Ref. [6] and has, as can be anticipated, about 1% smaller total uncertainty due to the inclusion of the exact NNLO for $q\bar{q} \rightarrow t\bar{t} + X$.

Example 8: The fixed order NNLO result at the Tevatron, with vanishing two-loop constant $C_{gg}^{(2)}$, is obtained from the file `top++Fixed-Order-TEV.cfg`:

$$\text{sigma_tot} = 7.00564 + 0.202329(2.88808\%) - 0.30846(4.40302\%) \text{ [pb]}.$$

This result corresponds to Eq. (11) in Ref. [5].

Appendix C. Program's structure: a brief overview

The function `main()` is located in the file `top++.cpp`. The program consists of four classes that are initialized in the function `main()`. The classes, listed in the order they are initialized, and their functionality are briefly described in the following.

1. **Class PartonicFlux.** This class precomputes the partonic fluxes on a grid of `NPdfGrid` points. One object of the class is created for each value of the factorization scale μ_F . See also Section 3.1.
2. **Class FixedOrder.** This class represents the partonic fixed order cross-section. One object of this class is created for each combination of (μ_F, μ_R) .
3. **Class SubtrFlux.** This class implements a fake partonic flux that mirrors the actual partonic flux within a distance $\mathcal{O}(\text{ETA})$ from the partonic threshold. Our implementation follows Ref. [15] where this flux was introduced and the need for it explained. Objects of class `SubtrFlux` are created for each object of class `PartonicFlux`.
4. **Class Resummation.** One object of this class is created for each pair (μ_F, μ_R) if a resummed calculation is requested by the user. Each object is constructed analytically in Mellin N -space and then inverted numerically back to x -space with the help of the Minimal Prescription of Ref. [15].

A small number of functions can be found outside the above classes. The file `lgamma` contains the logarithm of the Euler gamma function $\ln \Gamma(z)$ for complex argument, and `psin` contains the polygamma function $\Psi(z, k)$, $k = 0, 1$ for complex z . The remaining functions are located in the file `Utilities`. These are the prescriptions for calculation of PDF uncertainties (we have implemented two prescriptions; see the description of option `PDFuncertainty` in Section 2) and the actual computation of the final result as a convolution of perturbative functions and partonic fluxes.

- [1] M. Czakon and A. Mitov, arXiv:1207.0236 [hep-ph].
- [2] S. Alekhin, J. Blumlein and S. Moch, arXiv:1202.2281 [hep-ph].
- [3] S. Alekhin, J. Blumlein, S. Klein and S. Moch, Phys. Rev. D **81**, 014032 (2010) [arXiv:0908.2766 [hep-ph]].
- [4] S. Alekhin, Phys. Rev. D **68**, 014002 (2003) [hep-ph/0211096].
- [5] P. Baernreuther, M. Czakon and A. Mitov, arXiv:1204.5201 [hep-ph].
- [6] M. Cacciari, M. Czakon, M. L. Mangano, A. Mitov and P. Nason, Phys. Lett. B **710**, 612 (2012) [arXiv:1111.5869 [hep-ph]].
- [7] R. D. Ball *et al.* [The NNPDF Collaboration], Nucl. Phys. B **855**, 153 (2012) [arXiv:1107.2652 [hep-ph]].
- [8] P. M. Nadolsky and Z. Sullivan, eConfC **010630**, P510 (2001) [hep-ph/0110378].
- [9] M. Cacciari, S. Frixione, M. L. Mangano, P. Nason, G. Ridolfi, JHEP **0809**, 127 (2008). [arXiv:0804.2800 [hep-ph]].
- [10] A. D. Martin, W. J. Stirling, R. S. Thorne and G. Watt, Eur. Phys. J. C **63**, 189 (2009) [arXiv:0901.0002 [hep-ph]].
- [11] M. Beneke, M. Czakon, P. Falgari, A. Mitov, C. Schwinn, Phys. Lett. **B690**, 483-490 (2010). [arXiv:0911.5166 [hep-ph]].
- [12] R. Bonciani, S. Catani, M. L. Mangano, P. Nason, Nucl. Phys. **B529**, 424-450 (1998). [hep-ph/9801375].
- [13] M. R. Whalley, D. Bourilov and R. C. Group, arXiv:hep-ph/0508110; <http://hepforge.cedar.ac.uk/lhapdf/>.
- [14] M. Aliev, H. Lacker, U. Langenfeld, S. Moch, P. Uwer and M. Wiedermann, Comput. Phys. Commun. **182**, 1034 (2011) [arXiv:1007.1327 [hep-ph]].
- [15] S. Catani, M. L. Mangano, P. Nason and L. Trentadue, Nucl. Phys. B **478**, 273 (1996) [arXiv:hep-ph/9604351].